

# Extracting Transport Properties of Quark-Gluon Plasma from the Heavy-Quark Potential With Neural Networks in a Holographic Model

Wen-Chao Dai,<sup>1</sup> Ou-Yang Luo,<sup>1</sup> Bing Chen,<sup>1</sup> Xun Chen,<sup>1,2,\*</sup> Xiao-Yan Zhu,<sup>3,†</sup> and Xiao-Hua Li<sup>1,2,‡</sup>

<sup>1</sup>*School of Nuclear Science and Technology,  
University of South China, No. 28, West Changsheng Road,  
Hengyang City, Hunan Province, China.*

<sup>2</sup>*Key Laboratory of Advanced Nuclear Energy Design and Safety,  
Ministry of Education, Hengyang, 421001, China*

<sup>3</sup>*School of Mathematics and Physics,  
University of South China, No. 28, West Changsheng Road,  
Hengyang City, Hunan Province, China.*

(Dated: March 15, 2025)

## Abstract

Using Kolmogorov-Arnold Networks (KANs), we construct a holographic model informed by lattice QCD data. This neural network approach enables the derivation of an analytical solution for the deformation factor  $w(r)$  and the determination of a constant  $g$  related to the string tension. Within the KANs-based holographic framework, we further analyze heavy quark potentials under finite temperature and chemical potential conditions. Additionally, we calculate the drag force, jet quenching parameter, and diffusion coefficient of heavy quarks in this paper. Our findings demonstrate qualitative consistency with both experimental measurements and established phenomenological model.

---

\* [chenxunhep@qq.com](mailto:chenxunhep@qq.com)

† [xyzhu0128@163.com](mailto:xyzhu0128@163.com)

‡ [lixiaohuaphysics@126.com](mailto:lixiaohuaphysics@126.com)

## I. INTRODUCTION

High-energy nuclear collisions conducted at facilities such as the Relativistic Heavy-Ion Collider (RHIC) and the Large Hadron Collider (LHC) provide a unique opportunity to study the properties of the quark-gluon plasma (QGP), a deconfined state of matter formed under extreme temperatures and densities. The QGP offers a window into the fundamental aspects of QCD, and understanding its dynamics is crucial for advancing our knowledge of strong interactions. However, the computational and theoretical resolution of QCD still faces many challenges in practical applications and complex environments. To address these challenges, the Anti-de Sitter/conformal field theory (AdS/CFT) correspondence provides a way to solve the problem. Through this dual relationship, strong interactions can be explored in higher-dimensional spaces [1–3], offering new perspectives on the complex behavior of QCD.

Experimental and theoretical studies have found that heavy quarks are an important tool in the study of finite-temperature QCD matter due to their unique behavior under extreme conditions [4–17]. Under these conditions, the dissociation of heavy quark-antiquark pairs serves as a key indicator of strong interaction deconfinement, and thus it is critical to investigate the heavy quark potential within holographic QCD frameworks. Early work has tentatively revealed the characteristics of the holographic quark-antiquark pair potential energy [18–42]. However, as the study progresses, understanding the dynamics of the potentials between quarks becomes essential to reveal the complexity of QCD matter.

Jets and heavy quark diffusion are among the most useful probes for investigating the properties of the QGP at different scales. Jets, produced from high-energy partons in collisions, undergo complex processes such as energy loss, medium-induced radiation, and medium response as they traverse the QGP. The substructure of jets, particularly observables like the Energy-Energy Correlator (EEC), provides detailed insights into the interaction between jets and the QGP. The EEC is highly sensitive to the angular distribution of energy within jets, revealing the interplay of mass effects, energy loss, and medium response. Studies have shown a clear flavor hierarchy in the EEC for both vacuum and QGP environments, driven by the mass effect of heavy quarks (e.g., charm and bottom quarks). By analyzing heavy-flavor jets, researchers can probe the mass dependence of jet substructure and jet-medium interactions, offering a deeper understanding of QGP dynamics [43–50]. In parallel, the d-

iffusion of heavy quarks in the QGP is a critical aspect of understanding the hydrodynamic behavior of the plasma. Heavy quarks, such as charm and bottom quarks, are produced in the early stages of collisions and participate in the entire evolution of the QGP. The spatial diffusion coefficient, quantifies the momentum transfer from the QGP to heavy quarks and provides insights into the hydrodynamization process. Recent lattice QCD calculations with dynamical quarks have revealed that the heavy quark diffusion coefficient is significantly smaller than previous estimates from quenched lattice QCD and phenomenological models. This suggests that heavy quarks hydrodynamize very quickly in the QGP, highlighting the near-perfect fluidity of the medium [51–58].

In recent years, machine learning techniques, especially multilayer perceptrons (MLPs), have shown unprecedented potential in solving complex scientific problems [59–73]. According to the Universal Approximation Theorem, MLPs are able to approximate arbitrarily complex functions by increasing the number of neurons in the hidden layers, a feature that makes them perform well in solving partial differential equations (PDEs). Recently, Kolmogorov-Arnold Networks (KANs) have been proposed in Refs. [74, 75]. Unlike traditional MLPs, KANs fundamentally eliminate the reliance on linear weight matrices by using learnable functions instead of fixed activation functions. In terms of accuracy and interpretability, KANs even significantly outperform traditional MLPs for small-scale AI and scientific tasks, demonstrating superior performance and greater expressiveness [76, 77].

With the rapid development of machine learning technology [78–97], deep learning has received pioneering attention in the study of holographic QCD [98–100]. In addition, the combination of machine learning and holography has been deeply explored in a series of recent studies [101–112]. In contrast to traditional holographic models, this new approach first utilizes experimental or lattice QCD data to determine metrics and other model parameters with the help of machine learning. These acquired metrics are then applied to compute other physical QCD observables and are used as predictive outputs of the model. This interdisciplinary integration not only simplifies the complex computational process but also improves the prediction accuracy and reliability of the model, thereby opening up new directions in the study of strong interactions.

This paper is organized as follows: Section II will detail the calculation of the heavy quark potential using the Andreev-Zakharov model. In Section III, we construct a holographic model based on KANs and calculate the heavy-quark potential at finite temperatures and

chemical potentials. Our analysis will determine the critical temperature  $T_c$  and provide an analytical solution for  $w(r)$ . In Section IV, the drag and jet quenching parameters, as well as the diffusion coefficients of heavy quarks at finite temperature and finite chemical potential, will be calculated for both regimes based on the analytic form of the function  $w(r)$ . Finally, Section V will summarize the main results and conclusions of the whole paper.

## II. HOLOGRAPHIC HEAVY QUARK POTENTIAL

The Andreev-Zakharov model allows for an accurate description of the potential of heavy quarkonium [18, 20] and exotic hadron states [113–119] by introducing an ad hoc deformation factor for the AdS<sub>5</sub>-RN black hole. The background metric can be expressed as

$$ds^2 = w(r) \frac{1}{r^2} [-f(r)dt^2 + d\vec{x}^2 + f^{-1}(r)dr^2]. \quad (1)$$

$$f(r) = 1 - \left( \frac{1}{r_h^4} + q^2 r_h^2 \right) r^4 + q^2 r^6. \quad (2)$$

$q$  is the black hole charge,  $r_h$  is the position of the black hole horizon. The Hawking temperature of a black hole is defined as

$$T = \frac{1}{4\pi} \left| \frac{df}{dr} \right|_{r=r_h} = \frac{1}{\pi r_h} \left( 1 - \frac{1}{2} Q^2 \right), \quad (3)$$

where  $Q = qr_h^3$  and  $0 \leq Q \leq \sqrt{2}$ . The relationship between the chemical potential  $\mu$  and  $q$  is shown below

$$\mu = k \frac{Q}{r_h}. \quad (4)$$

In this paper, we fix the parameter  $k$  to 1. Thus, we get

$$f(r) = 1 - \left( \frac{1}{r_h^4} + \frac{\mu^2}{r_h^2} \right) r^4 + \frac{\mu^2}{r_h^4} r^6. \quad (5)$$

$$T = \frac{1}{\pi r_h} \left( 1 - \frac{1}{2} \mu^2 r_h^2 \right). \quad (6)$$

If we choose the static norms  $t = \tau$  and  $\sigma = x$ , then a static quark-antiquark pair locating at

$$x = -\frac{L}{2}, \text{ and } x = \frac{L}{2}. \quad (7)$$

The Nambu-Goto action of a  $U$ -type string can be expressed as

$$S = \frac{1}{2\pi\alpha'} \int d\tau d\sigma \sqrt{-\det(g_{\alpha\beta})}, \quad (8)$$

with

$$g_{\alpha\beta} = G_{\mu\nu} \frac{\partial x^\mu}{\partial \sigma^\alpha} \frac{\partial x^\nu}{\partial \sigma^\beta}. \quad (9)$$

Therefore, it can be

$$S = \frac{g}{2\pi T} \int_{-\frac{L}{2}}^{\frac{L}{2}} dx \frac{w(r)}{r^2} \sqrt{f(r) + (\partial_x r)^2}, \quad (10)$$

where,  $g = \frac{1}{\alpha'}$  is related to the string tension,  $\alpha'$  is the square of the string length parameter in string theory. We now define the Lagrangian quantity as

$$\mathcal{L} = \frac{w(r)}{r^2} \sqrt{f(r) + (\partial_x r)^2}. \quad (11)$$

Then we have conserved quantity

$$\mathcal{H} = \frac{w(r)f(r)}{r^2 \sqrt{f(r) + (\partial_x r)^2}}. \quad (12)$$

At the maximum position of the U-shape string  $r_0$ , we have

$$\frac{w(r)f(r)}{r^2 \sqrt{f(r) + (\partial_x r)^2}} = \frac{w(r_0)}{r_0^2} \sqrt{f(r_0)}. \quad (13)$$

As mentioned before,  $\partial_x r$  can be solved as

$$\partial_x r = \sqrt{\frac{w^2(r)f^2(r)/r^4 - w^2(r_0)f(r_0)f(r)/r_0^4}{w^2(r_0)f(r_0)/r_0^4}}. \quad (14)$$

Therefore,  $\partial_r x$  is

$$\partial_r x = \sqrt{\frac{w^2(r_0)f(r_0)/r_0^4}{w^2(r)f^2(r)/r^4 - w^2(r_0)f(r_0)f(r)/r_0^4}}. \quad (15)$$

The quark distance  $L$  is defined as

$$L = 2 \int_0^{r_0} \partial_r x dr = 2 \int_0^{r_0} \sqrt{\frac{f^2(r_0)w^2(r_0)/r_0^4}{f^2(r)f(r_0)w^2(r)/r^4 - f^2(r_0)f(r)w^2(r_0)/r_0^4}} dr. \quad (16)$$

The heavy-quark potential can be written as

$$\begin{aligned} E &= 2g \int_0^{r_0} \left( \frac{w(r)}{r^2} \sqrt{1 + f(r)(\partial_r x)^2} - \frac{w(0)}{r^2} - \frac{w'(0)}{r} \right) dr - 2\frac{g}{r_0} w(0) + 2gw'(0)\ln(r_0) \\ &= 2g \int_0^{r_0} \left( \frac{w(r)}{r^2} \sqrt{1 + \frac{f(r_0)w^2(r_0)/r_0^4}{f(r)w^2(r)/r^4 - f(r_0)w^2(r_0)/r_0^4}} - \frac{w(0)}{r^2} - \frac{w'(0)}{r} \right) dr \\ &\quad - 2\frac{g}{r_0} w(0) + 2gw'(0)\ln(r_0). \end{aligned} \quad (17)$$

In the Andreev-Zakharov model, the string tension is  $g = 0.176$ , and the function  $w(r)$  is given by  $w(r) = e^{sr^2}$ , where  $s = 0.45$ . At zero temperature, we only need to set  $f(r) = 1$ . In the next section, we will first use the KANs to construct the holographic model at vanishing temperature.

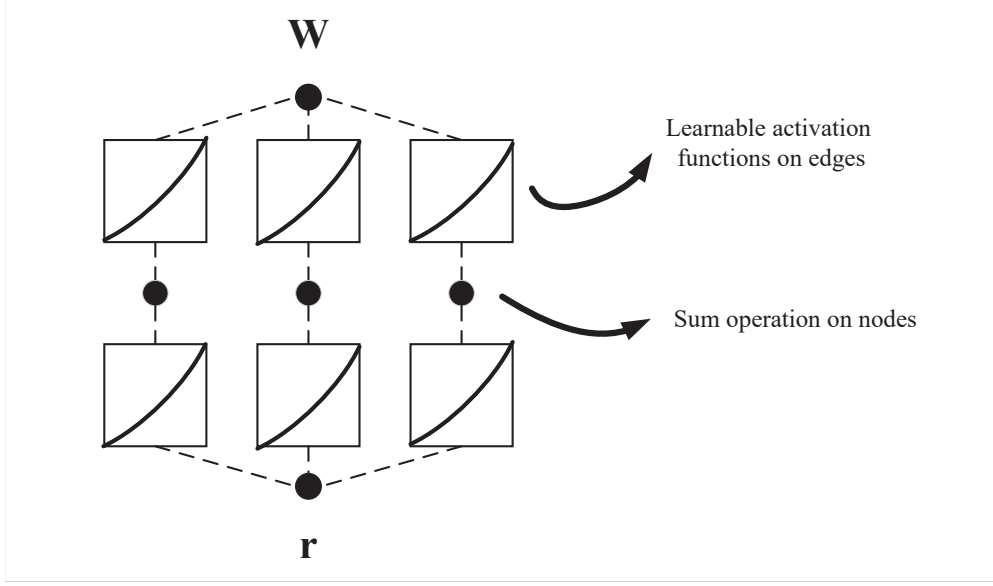


FIG. 1. A sketch of the KANs architecture used in our paper.

### III. CONSTRUCTION OF A KANS-BASED HOLOGRAPHIC MODEL

Based on KANs, we reconstruct the deformation factor  $w(r)$  in this section. We assume that  $w(r)$  is a specific function derived from the lattice data, with the string tension  $g$  used as a free parameter. The structure of the KANs we designed is shown in Fig. 1, we performed a reconstruction with the string tension  $g = 0.2425$  for the 2-flavor. The trained function  $w(r)$  is  $w(r) = 5.32 - 4.53 \sin \left( 11.16 - 3.11e^{-0.04(0.93r-1)^2} \right)$ , which is substituted into Eq. (17) for potential energy calculations. These calculations are combined with the lattice QCD data [120], and the results are shown in Fig. 2.

From Fig. 2 (a), it can be clearly observed that the reconstructed function satisfies the boundary condition  $w(0) \rightarrow 1$  and the function  $w(r)$  is increasing with  $r$ . Fig. 2 (b) illustrates the fitting performance of the neural network to the training dataset. These curves clearly demonstrate the excellent ability of the network in modeling the function  $E(L)$  with an accuracy that highly matches the theoretical expectation. This result further supports the validity of the model, verifies the reliability of the KANs.

In the subsequent stage of this study, we obtain an analytical solution of  $w(r)$ . From this solution, we calculate the heavy quark potential under finite temperature and chemical potential conditions by adding the function  $f(r)$ . Fig. 3 (a) shows that the linear component of the potential energy exhibits a decaying trend under finite temperature conditions.

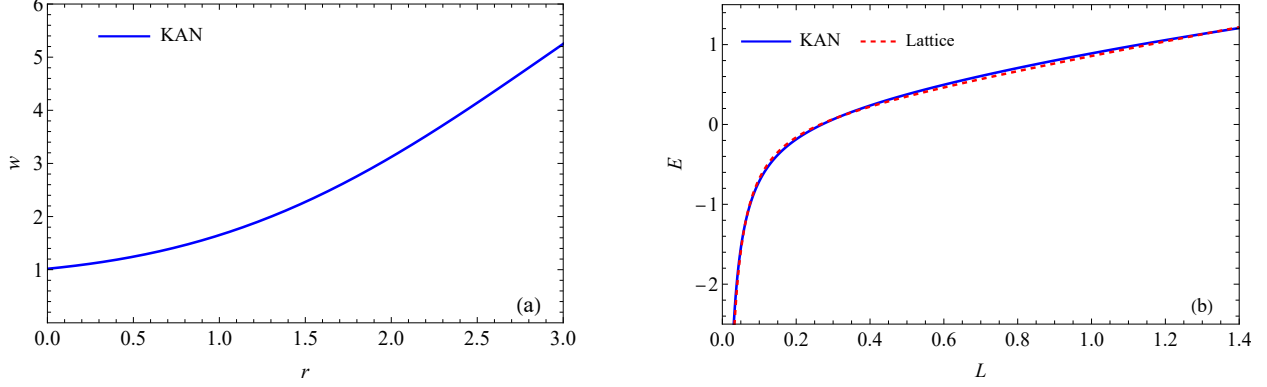


FIG. 2. (a) The deformation factor  $w(r)$  trained by KANs. (b) Comparison of the potential energy for the 2-flavor lattice with KANs. The blue solid line represents the calculation results of KANs, and the red dashed line represents the 2-flavor lattice data [120].

This decay may be attributed to the temperature-induced shielding effect. In contrast, the Coulombic component of the potential energy shows remarkable robustness to temperature changes and is almost unaffected by temperature increase. As the temperature increases, the coupling strength of the strong interactions weakens, leading to weaker inter-quark binding and the potential energy tends to vanish on a smaller spatial scale. This behavior is consistent with the results of the literature [121], further verifying the reliability of the analytic solution of  $w(r)$  obtained through neural networks. Fig. 3 (b) shows a similar trend. However, by comparison (a), we find that the chemical potential affects the potential energy to a significantly lesser extent than temperature does. These observations reveal the differential effects of temperature and chemical potential on the strong interaction potential energy.

At larger distance scales, the interaction between quarks and antiquarks is significantly weakened so that they behave as free particles. This phenomenon suggests that when the separation distance between quarks and antiquarks is large enough, they can be regarded as independently existing states and are no longer bound by strong interactions. The expectation value of the Polyakov loop can be defined as [20, 22]

$$\langle P \rangle = \exp \left\{ -\frac{1}{2T} E(r = \infty, T) \right\}. \quad (18)$$

With this choice, the Polyakov loop expectation value takes the form shown in Fig 4 (a). From Fig. 4 (b), it can be clearly seen that at  $T_c = 0.17$  GeV, the slope of the Polyakov loop expectation value is the largest.

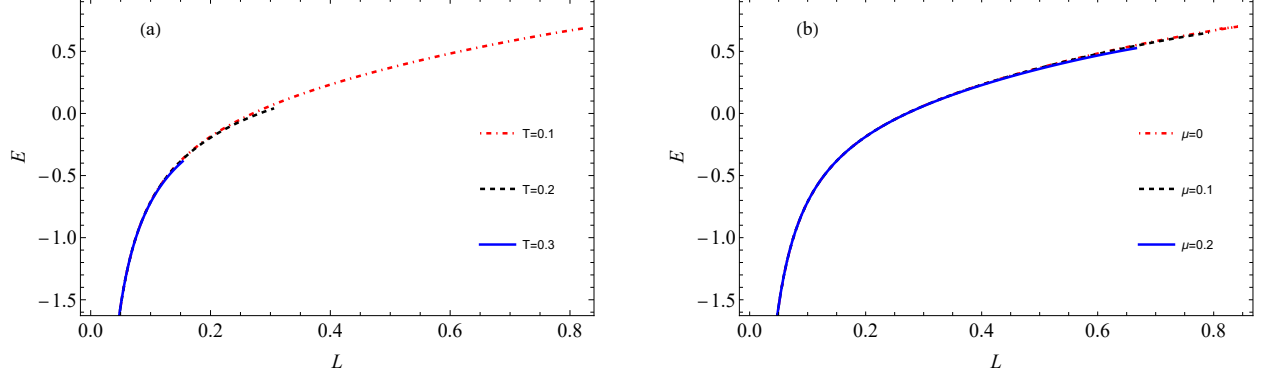


FIG. 3. (a) Potential energy  $E$  of quark-antiquark pairs as a function of quark separation distance  $L$  at different temperatures  $T$  when chemical potential  $\mu=0$ . (b) Variation of quark separation distance  $L$  with potential energy  $E$  for different chemical potentials at fixed temperature  $T = 0.1$ . The unit of  $E$  is GeV, the unit of  $L$  is fm, the unit of  $\mu$  is GeV, the unit of  $T$  is GeV.

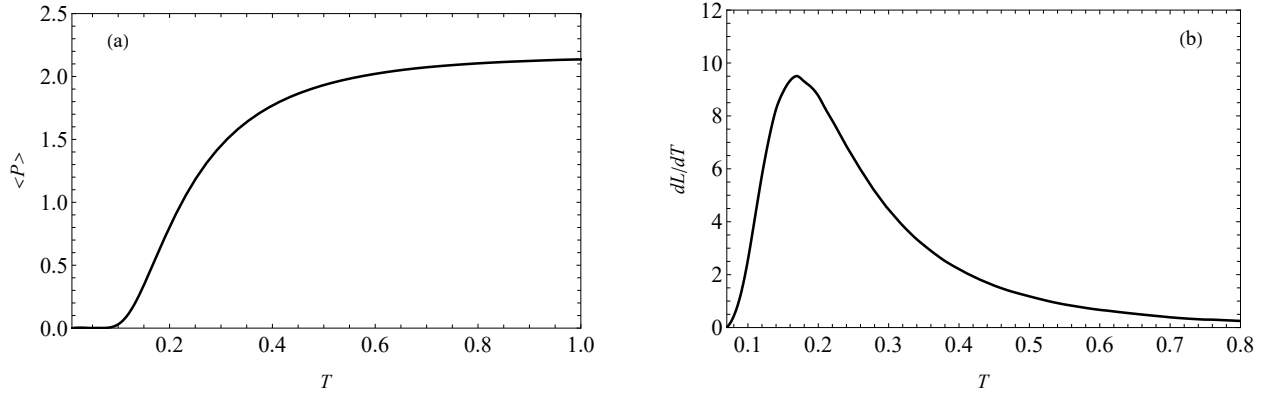


FIG. 4. (a) The relationship between the Polyakov loop expectation value  $\langle P \rangle$  and temperature  $T$ . (b) The relationship between the slope of the Polyakov loop expectation value  $\frac{d\langle P \rangle}{dT}$  and temperature  $T$ . The unit of  $T$  is GeV.

#### IV. TRANSPORT PROPERTIES OF QGP

In this section, we extend the calculation of analytic solutions of the function  $w(r)$  with string tension  $g$  to the drag force, diffusion coefficient, and jet quenching parameter of heavy quarks at finite temperature and finite chemical potential. To facilitate the study of holographic probes, we define  $A_s(r) = \frac{1}{2} \log(w(r))$ . Following Ref. [122], the drag force can be obtained by:

$$F_{\text{drag}} = \frac{dp}{dt} = \frac{dE}{dx} = -\frac{1}{2\pi\alpha'} \frac{e^{2A_s(r_s)} v}{r_s^2}, \quad (19)$$



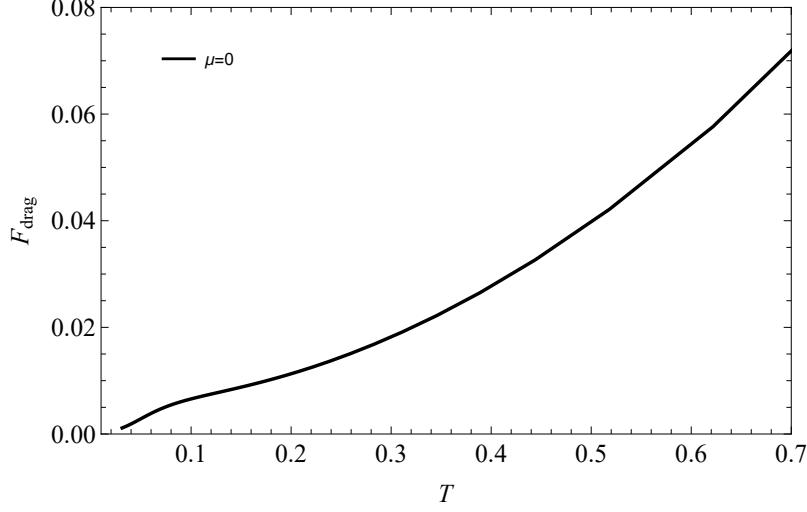


FIG. 5. Variation of drag force with temperature when quarks are in  $v = 0.3$  and  $\mu = 0$ .  $T$  and  $\mu$  are in units of GeV.

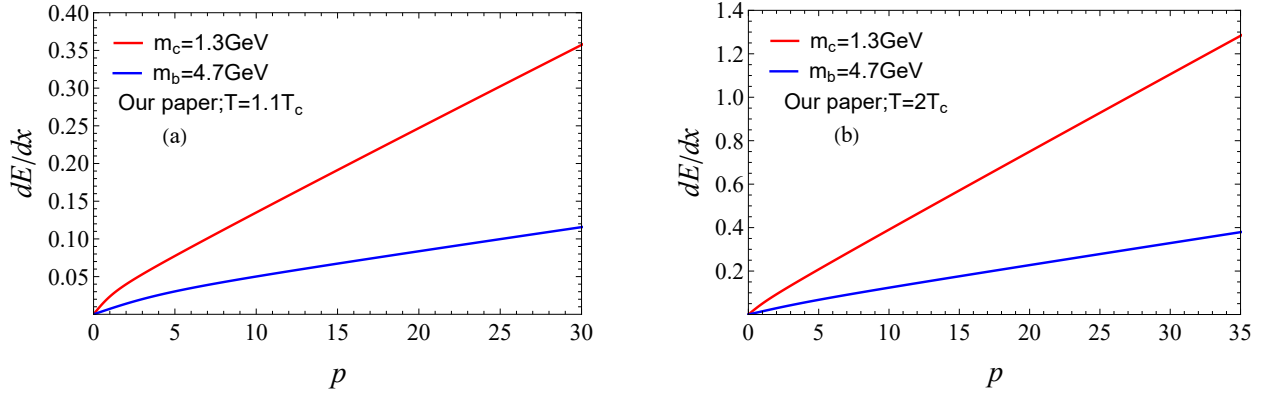


FIG. 6. Variation of energy loss with momentum  $p$  for bottom ( $m_b = 4.7$  GeV) and charm quark ( $m_c = 1.3$  GeV). (a) Variation of energy loss with momentum  $p$  at temperature  $T = 1.1T_c$ . (b) Energy loss versus momentum at temperature  $T = 2T_c$ .

where  $r_s$  satisfies  $f(r_s) - v^2 = 0$ . In this context, we first need to solve Eq. (5) numerically to obtain  $r$ , and subsequently employ Eq. (18) to compute the drag force. In Fig. 5, we show the variation of drag force with temperature at vanishing chemical potential. As can be seen from the figure, the drag force increases significantly with increasing temperature.

According to Eq. (19) an be deduced that the energy loss is equal to the drag force, which allows us to plot the relationship between energy loss and momentum in different systems. We have

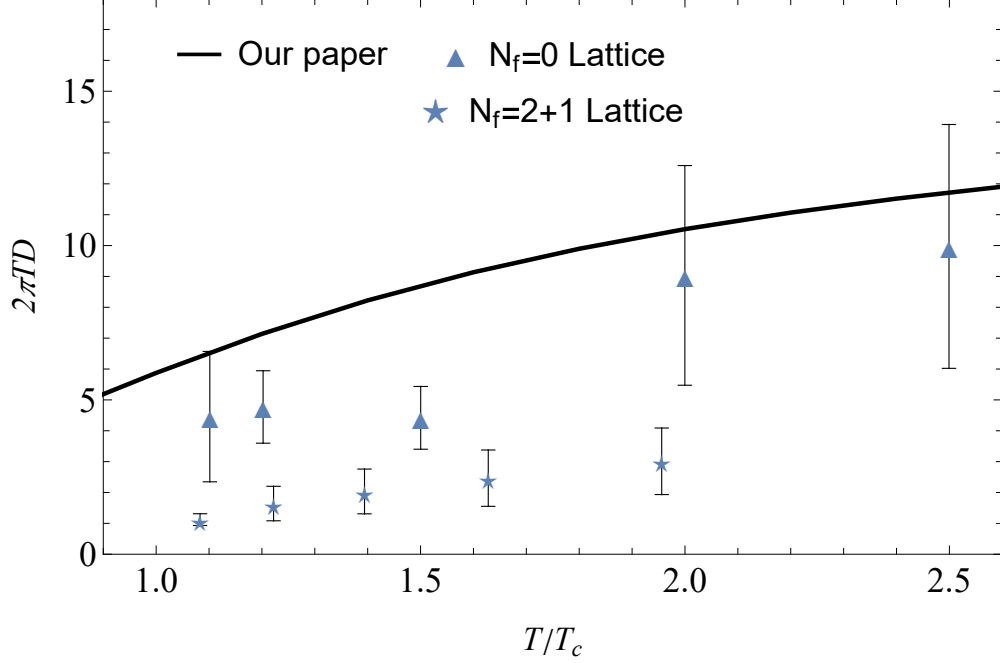


FIG. 7. The figure shows the normalized diffusion coefficient versus  $T/T_c$ . The unit of  $T$  is GeV.

$$\frac{dE}{dx} = -\frac{e^{2A_s(r_s)}\sqrt{1-v^2}\pi T^2\sqrt{\lambda}}{\pi^2 T^2 r_s^2} \frac{\pi T^2\sqrt{\lambda}}{2m} p. \quad (20)$$

Fig. 6 illustrates the energy loss of bottom quark ( $m_b = 4.7$  GeV) and charm quark ( $m_c = 1.3$  GeV) at vanishing chemical potential. From Fig. 6, it is clear that the energy loss increases with momentum. In addition, higher temperatures lead to an increase in energy loss. The qualitative results are similar to Refs. [122–124].

Next we proceed to study the diffusion coefficient, which in the AdS/Schwarzschild context can be written as [122]

$$D = \frac{T}{m} t = \frac{2}{\pi T g} \frac{\pi^2 T^2 r_s^2}{e^{2A_s(r_s)}\sqrt{1-v^2}}. \quad (21)$$

According to Eq. (21), we calculate the diffusion coefficients of space heavy quarks normalized by  $2\pi T$ , as shown in Fig. 7. As can be seen from the figure, the diffusion coefficient of heavy quarks gradually increases with increasing temperature, and this behavior is consistent with reference [51], indicating that the analytic solution of  $w(r)$  obtained by KANs is reliable. Besides, the qualitative behavior of diffusion coefficient is consistent with Ref. [125].

Now we turn to the study of jet quenching parameter and we obtain the following ex-

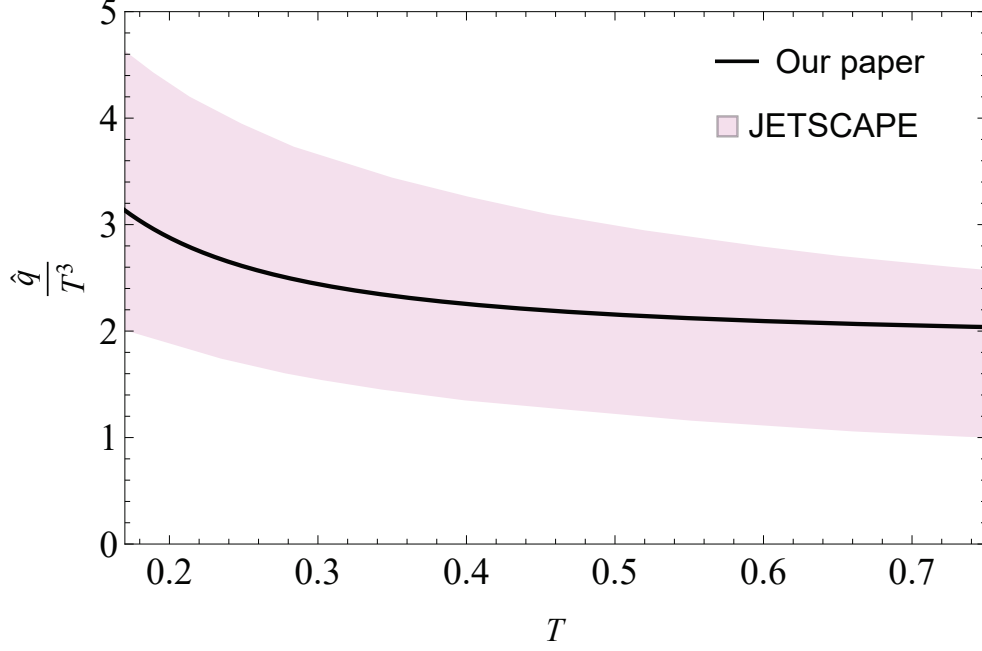


FIG. 8. Relationship between  $\hat{q}/T^3$  and temperature  $T$  at zero chemical potential. The unit of  $T$  is GeV.

pression for the jet quenching parameter in the holographic model [122]

$$\hat{q} = \frac{g}{\pi a_0}. \quad (22)$$

Here  $a_0$  is defined

$$a_0 = \int_0^{r_h} \frac{dr r^2 L^{-2} e^{-2A_s(r)}}{\sqrt{f(r)(1-f(r))}}. \quad (23)$$

To obtain the jet quenching parameters in the holographic QCD model, we numerically solved using Eq. (22), comparing  $\hat{q}/T^3$  as a function of temperature, as shown in Fig. 8. From Fig. 9, it can be observed that temperature leads to an enhancement of the jet quenching parameter, indicating that in the considered model, a denser or hotter medium results in increased energy loss, which aligns with the physical intuition that jets passing through a higher-temperature (i.e., higher-density or more particle-rich) medium encounter more scattering centers and therefore experience greater energy loss. The results of our model calculations are consistent with the experimental results of RHIC and LHC [47].

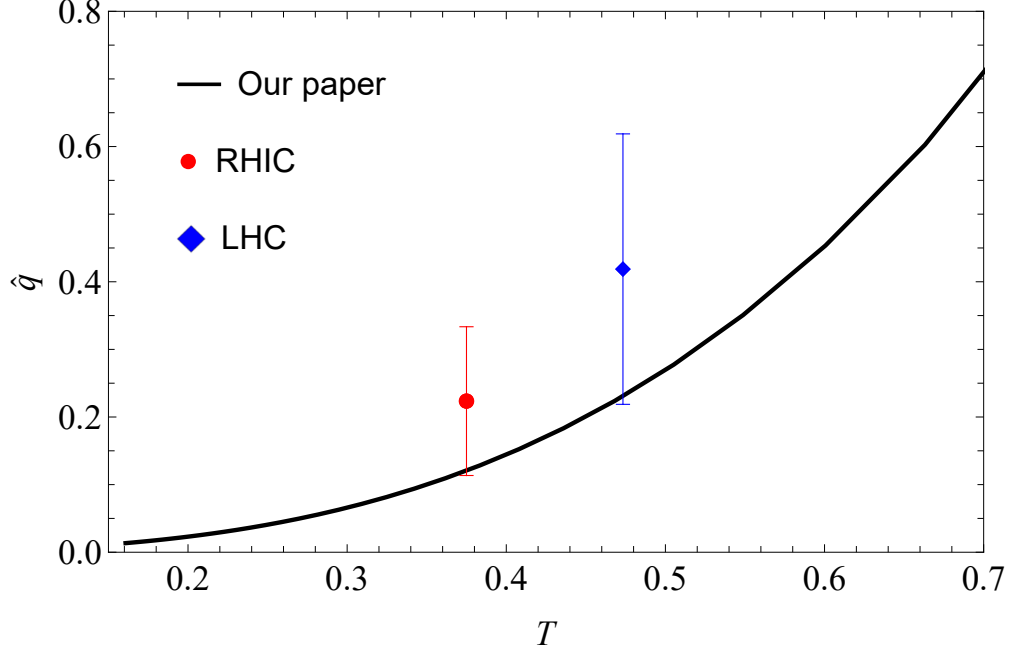


FIG. 9. The relationship between the jet quenching parameter and temperature in the holographic model at zero chemical potential. Error bars represent experimental values from RHIC and LHC [47]. The unit of  $T$  is GeV.

## V. SUMMARY

In this study, we employ the KANs to extract data from QCD to construct a holographic models. In order to verify the validity of the reconstruction results, we first apply the obtained function  $w(r)$  to the heavy-quark potential and compare with lattice data. The results show that KANs exhibit significant effectiveness in solving inverse problems. It is worth emphasizing that KANs have the ability to provide analytical solutions. Moreover, we further examine the heavy quark potential and its relationship at finite temperature and chemical potential.

In addition, based on the constructed function  $w(r)$ , we study the relationship between the drag force of heavy quarks, the diffusion coefficient, and the jet quenching parameter under the conditions of finite temperature and finite chemical potential, which reveals the accuracy of the KANs-based holographic model. This finding lays the foundation for understanding the complementarity of different models in dealing with complex physical phenomena, and also provides new perspectives for future research.

The findings not only provide an effective paradigm for utilizing machine learning methods

to solve complex physics problems, but also point to new directions for subsequent research. These new directions include an in-depth exploration of the interactions between different physical fields, and the use of holographic models in combination with machine learning methods, with a view to a more comprehensive understanding of the dynamic behavior of complex systems.

Finally, we want to emphasize that our results focus on qualitative behavior. On the one hand, we can develop a holographic model derived from the Einstein equation with KANs. On the other hand, we should input more information into the holographic model so that it can provide precise predictions of the observables in the future.

## ACKNOWLEDGEMENT

This work is supported by National Natural Science Foundation of China (NSFC) Grants No. 12405154, 12175100.

## REFERENCES

- 
- [1] J. M. Maldacena, *Adv. Theor. Math. Phys.* **2**, 231 (1998), [arXiv:hep-th/9711200](#).
  - [2] E. Witten, *Adv. Theor. Math. Phys.* **2**, 505 (1998), [arXiv:hep-th/9803131](#).
  - [3] S. S. Gubser, I. R. Klebanov, and A. M. Polyakov, *Phys. Lett. B* **428**, 105 (1998), [arXiv:hep-th/9802109](#).
  - [4] A. Andronic *et al.*, *Eur. Phys. J. C* **76**, 107 (2016), [arXiv:1506.03981 \[nucl-ex\]](#).
  - [5] X. Luo and N. Xu, *Nucl. Sci. Tech.* **28**, 112 (2017), [arXiv:1701.02105 \[nucl-ex\]](#).
  - [6] Z. Tang, Z.-B. Tang, W. Zha, W.-M. Zha, Y. Zhang, and Y.-F. Zhang, *Nucl. Sci. Tech.* **31**, 81 (2020), [arXiv:2105.11656 \[nucl-ex\]](#).
  - [7] F. Liu, *Nucl. Sci. Tech.* **24**, 50524 (2013).
  - [8] K. Hattori and X.-G. Huang, *Nucl. Sci. Tech.* **28**, 26 (2017), [arXiv:1609.00747 \[nucl-th\]](#).
  - [9] Q.-Y. Shou *et al.*, *Nucl. Sci. Tech.* **35**, 219 (2024), [arXiv:2409.17964 \[nucl-ex\]](#).
  - [10] C. M. Ko and F. Li, *Nucl. Sci. Tech.* **27**, 140 (2016).
  - [11] R. Averbeck, *Prog. Part. Nucl. Phys.* **70**, 159 (2013), [arXiv:1505.03828 \[nucl-ex\]](#).

- [12] G. Aarts, C. Allton, M. B. Oktay, M. Peardon, and J.-I. Skullerud, *Phys. Rev. D* **76**, 094513 (2007), [arXiv:0705.2198 \[hep-lat\]](#).
- [13] T. Matsui and H. Satz, *Phys. Lett. B* **178**, 416 (1986).
- [14] N. Brambilla *et al.* (Quarkonium Working Group), (2004), [10.5170/CERN-2005-005](#), [arXiv:hep-ph/0412158](#).
- [15] R. Rapp and H. van Hees (2010) pp. 111–206, [arXiv:0903.1096 \[hep-ph\]](#).
- [16] S. Zheng, J. Liu, S. Zheng, J. Zhao, and B. Chen, (2024), [arXiv:2411.12169 \[nucl-th\]](#).
- [17] G. Chen, B. Chen, and J. Zhao, *Eur. Phys. J. C* **84**, 869 (2024), [arXiv:2402.11316 \[nucl-th\]](#).
- [18] O. Andreev and V. I. Zakharov, *Phys. Rev. D* **74**, 025023 (2006), [arXiv:hep-ph/0604204](#).
- [19] O. Andreev and V. I. Zakharov, *Phys. Lett. B* **645**, 437 (2007), [arXiv:hep-ph/0607026](#).
- [20] O. Andreev and V. I. Zakharov, *JHEP* **04**, 100 (2007), [arXiv:hep-ph/0611304](#).
- [21] S. He, M. Huang, and Q.-s. Yan, *Prog. Theor. Phys. Suppl.* **186**, 504 (2010), [arXiv:1007.0088 \[hep-ph\]](#).
- [22] P. Colangelo, F. Giannuzzi, and S. Nicotri, *Phys. Rev. D* **83**, 035015 (2011), [arXiv:1008.3116 \[hep-ph\]](#).
- [23] O. DeWolfe, S. S. Gubser, and C. Rosen, *Phys. Rev. D* **83**, 086005 (2011), [arXiv:1012.1864 \[hep-th\]](#).
- [24] D. Li, S. He, M. Huang, and Q.-S. Yan, *JHEP* **09**, 041 (2011), [arXiv:1103.5389 \[hep-th\]](#).
- [25] K. B. Fadafan, *Eur. Phys. J. C* **71**, 1799 (2011), [arXiv:1102.2289 \[hep-th\]](#).
- [26] K. B. Fadafan and E. Azimfard, *Nucl. Phys. B* **863**, 347 (2012), [arXiv:1203.3942 \[hep-th\]](#).
- [27] R.-G. Cai, S. He, and D. Li, *JHEP* **03**, 033 (2012), [arXiv:1201.0820 \[hep-th\]](#).
- [28] D. Li, M. Huang, and Q.-S. Yan, *Eur. Phys. J. C* **73**, 2615 (2013), [arXiv:1206.2824 \[hep-th\]](#).
- [29] Z. Fang, S. He, and D. Li, *Nucl. Phys. B* **907**, 187 (2016), [arXiv:1512.04062 \[hep-ph\]](#).
- [30] Y. Yang and P.-H. Yuan, *JHEP* **12**, 161 (2015), [arXiv:1506.05930 \[hep-th\]](#).
- [31] Z.-q. Zhang, D.-f. Hou, and G. Chen, *Nucl. Phys. A* **960**, 1 (2017), [arXiv:1507.07263 \[hep-ph\]](#).
- [32] C. Ewerz, O. Kaczmarek, and A. Samberg, *JHEP* **03**, 088 (2018), [arXiv:1605.07181 \[hep-th\]](#).
- [33] X. Chen, S.-Q. Feng, Y.-F. Shi, and Y. Zhong, *Phys. Rev. D* **97**, 066015 (2018), [arXiv:1710.00465 \[hep-ph\]](#).
- [34] I. Aref’eva and K. Rannu, *JHEP* **05**, 206 (2018), [arXiv:1802.05652 \[hep-th\]](#).

- [35] H. Bohra, D. Dudal, A. Hajilou, and S. Mahapatra, *Phys. Lett. B* **801**, 135184 (2020), [arXiv:1907.01852 \[hep-th\]](#).
- [36] X. Chen, D. Li, D. Hou, and M. Huang, *JHEP* **03**, 073 (2020), [arXiv:1908.02000 \[hep-ph\]](#).
- [37] J. Zhou, X. Chen, Y.-Q. Zhao, and J. Ping, *Phys. Rev. D* **102**, 086020 (2020), [arXiv:2006.09062 \[hep-ph\]](#).
- [38] J. Zhou, X. Chen, Y.-Q. Zhao, and J. Ping, *Phys. Rev. D* **102**, 126029 (2021).
- [39] X. Chen, L. Zhang, D. Li, D. Hou, and M. Huang, *JHEP* **07**, 132 (2021), [arXiv:2010.14478 \[hep-ph\]](#).
- [40] X. Chen, L. Zhang, and D. Hou, *Chin. Phys. C* **46**, 073101 (2022), [arXiv:2108.03840 \[hep-ph\]](#).
- [41] X. Guo, X. Chen, D. Xiang, M. A. Martin Contreras, and X.-H. Li, *Phys. Rev. D* **110**, 046014 (2024), [arXiv:2406.04650 \[hep-ph\]](#).
- [42] L. Bellantuono, P. Colangelo, F. De Fazio, F. Giannuzzi, and S. Nicotri, *Nucl. Phys. A* **982**, 931 (2019).
- [43] W.-J. Xing, S. Cao, G.-Y. Qin, and X.-N. Wang, *Phys. Rev. Lett.* **134**, 052301 (2025), [arXiv:2409.12843 \[hep-ph\]](#).
- [44] W.-J. Xing, S. Cao, and G.-Y. Qin, *Phys. Lett. B* **850**, 138523 (2024), [arXiv:2303.12485 \[hep-ph\]](#).
- [45] S. Cao *et al.* (JETSCAPE), *Phys. Rev. C* **104**, 024905 (2021), [arXiv:2102.11337 \[nucl-th\]](#).
- [46] M. Xie, W. Ke, H. Zhang, and X.-N. Wang, *Phys. Rev. C* **108**, L011901 (2023), [arXiv:2206.01340 \[hep-ph\]](#).
- [47] K. M. Burke *et al.* (JET), *Phys. Rev. C* **90**, 014909 (2014), [arXiv:1312.5003 \[nucl-th\]](#).
- [48] A. Majumder and M. Van Leeuwen, *Prog. Part. Nucl. Phys.* **66**, 41 (2011), [arXiv:1002.2206 \[hep-ph\]](#).
- [49] S. Wang, S. Li, Y. Li, B.-W. Zhang, and E. Wang, (2024), [10.1088/1674-1137/adb385](#), [arXiv:2410.21834 \[hep-ph\]](#).
- [50] X.-Q. Cai, S. Li, Z.-F. Liu, and T.-Y. Wang, *Sci. Rep.* **14**, 18385 (2024).
- [51] L. Altenkort, O. Kaczmarek, R. Larsen, S. Mukherjee, P. Petreczky, H.-T. Shu, and S. Stenkebach (HotQCD), *Phys. Rev. Lett.* **130**, 231902 (2023), [arXiv:2302.08501 \[hep-lat\]](#).
- [52] J. Casalderrey-Solana and D. Teaney, *Phys. Rev. D* **74**, 085012 (2006), [arXiv:hep-ph/0605199](#).

- [53] S. Caron-Huot and G. D. Moore, *Phys. Rev. Lett.* **100**, 052301 (2008), [arXiv:0708.4232 \[hep-ph\]](#).
- [54] L. Altenkort, A. M. Eller, O. Kaczmarek, L. Mazur, G. D. Moore, and H.-T. Shu, *Phys. Rev. D* **103**, 014511 (2021), [arXiv:2009.13553 \[hep-lat\]](#).
- [55] N. Brambilla, V. Leino, P. Petreczky, and A. Vairo, *Phys. Rev. D* **102**, 074503 (2020), [arXiv:2007.10078 \[hep-lat\]](#).
- [56] D. Banerjee, S. Datta, R. V. Gavai, and P. Majumdar, *Nucl. Phys. A* **1038**, 122721 (2023), [arXiv:2206.15471 \[hep-ph\]](#).
- [57] S. Acharya *et al.* (ALICE), *JHEP* **01**, 174 (2022), [arXiv:2110.09420 \[nucl-ex\]](#).
- [58] Z.-R. Zhu, S.-Q. Feng, Y.-F. Shi, and Y. Zhong, *Phys. Rev. D* **99**, 126001 (2019), [arXiv:1901.09304 \[hep-ph\]](#).
- [59] Y.-G. Ma, L.-G. Pang, R. Wang, and K. Zhou, *Chin. Phys. Lett.* **40**, 122101 (2023), [arXiv:2311.07274 \[nucl-th\]](#).
- [60] K. Zhou, L. Pang, S. Shi, and H. Stoecker, *PoS FAIRness2022*, 064 (2023).
- [61] L.-G. Pang, *Int. J. Mod. Phys. E* **33**, 2430009 (2024).
- [62] Z.-H. Li, C.-Q. Li, and L.-G. Pang, (2023), [arXiv:2309.07397 \[gr-qc\]](#).
- [63] K.-F. Pu, H. Li, H.-L. Lu, and L.-G. Pang, *Chin. Phys. C* **47**, 054104 (2023), [arXiv:2303.03934 \[nucl-th\]](#).
- [64] L.-G. Pang, *Nucl. Phys. A* **1005**, 121972 (2021).
- [65] J. Steinheimer, L. Pang, K. Zhou, V. Koch, J. Randrup, and H. Stoecker, *JHEP* **12**, 122 (2019), [arXiv:1906.06562 \[nucl-th\]](#).
- [66] L. Wang, B. M. Hare, K. Zhou, H. Stöcker, and O. Scholten, *Chaos Solitons Fractals* **170**, 113346 (2023).
- [67] Y. Huang, J. Chen, J. Jia, L.-M. Liu, Y.-G. Ma, and C. Zhang, (2025), [arXiv:2501.01352 \[nucl-th\]](#).
- [68] W.-B. He, Y.-G. Ma, L.-G. Pang, H.-C. Song, and K. Zhou, *Nucl. Sci. Tech.* **34**, 88 (2023), [arXiv:2303.06752 \[hep-ph\]](#).
- [69] J. He, W.-B. He, Y.-G. Ma, and S. Zhang, *Phys. Rev. C* **104**, 044902 (2021), [arXiv:2109.06277 \[hep-ph\]](#).
- [70] R. Wang, Y.-G. Ma, R. Wada, L.-W. Chen, W.-B. He, H.-L. Liu, and K.-J. Sun, *Phys. Rev. Res.* **2**, 043202 (2020), [arXiv:2010.15043 \[nucl-th\]](#).



- [71] R. Guo, Y. Li, and B. Chen, [Entropy](#) **25**, 1563 (2023), [arXiv:2311.02335 \[nucl-th\]](#).
- [72] J.-G. Deng, J.-X. Li, J.-H. Cheng, L. Mu, and H. Zhang, [Chinese Physics C](#) (2025).
- [73] W. Kou and X. Chen, (2024), [arXiv:2411.14902 \[hep-ph\]](#).
- [74] Z. Liu, Y. Wang, S. Vaidya, F. Ruehle, J. Halverson, M. Soljačić, T. Y. Hou, and M. Tegmark, (2024), [arXiv:2404.19756 \[cs.LG\]](#).
- [75] Z. Liu, P. Ma, Y. Wang, W. Matusik, and M. Tegmark, “[Kan 2.0: Kolmogorov-arnold networks meet science](#),” (2024), [arXiv:2408.10205 \[cs.LG\]](#).
- [76] J. Erdmann, F. Mausolf, and J. L. Späh, (2024), [arXiv:2408.02743 \[hep-ph\]](#).
- [77] L. Yu *et al.*, (2024), [arXiv:2412.18301 \[astro-ph.IM\]](#).
- [78] W. He, Q. Li, Y. Ma, Z. Niu, J. Pei, and Y. Zhang, [Sci. China Phys. Mech. Astron.](#) **66**, 282001 (2023), [arXiv:2301.06396 \[nucl-th\]](#).
- [79] K. Zhou, L.-G. Pang, N. Su, H. Petersen, H. Stoecker, and X.-N. Wang, [EPJ Web Conf.](#) **171**, 16005 (2018).
- [80] L.-G. Pang, K. Zhou, N. Su, H. Petersen, H. Stöcker, and X.-N. Wang, [Nucl. Phys. A](#) **982**, 867 (2019).
- [81] L. Wang, Y. Jiang, L. He, and K. Zhou, [Chin. Phys. Lett.](#) **39**, 120502 (2022), [arXiv:2005.04857 \[cond-mat.dis-nn\]](#).
- [82] Y.-S. Zhao, L. Wang, K. Zhou, and X.-G. Huang, [Phys. Rev. C](#) **106**, L051901 (2022), [arXiv:2105.13761 \[hep-ph\]](#).
- [83] Y.-L. Du, D. Pablos, and K. Tywoniuk, [JHEP](#) **21**, 206 (2020), [arXiv:2012.07797 \[hep-ph\]](#).
- [84] Y.-L. Du, D. Pablos, and K. Tywoniuk, [Phys. Rev. Lett.](#) **128**, 012301 (2022), [arXiv:2106.11271 \[hep-ph\]](#).
- [85] L. Jiang, L. Wang, and K. Zhou, [Phys. Rev. D](#) **103**, 116023 (2021), [arXiv:2103.04090 \[nucl-th\]](#).
- [86] S. Shi, K. Zhou, J. Zhao, S. Mukherjee, and P. Zhuang, [PoS LATTICE2021](#), 537 (2022).
- [87] S. Soma, L. Wang, S. Shi, H. Stöcker, and K. Zhou, [JCAP](#) **08**, 071 (2022), [arXiv:2201.01756 \[hep-ph\]](#).
- [88] F.-P. Li, H.-L. Lü, L.-G. Pang, and G.-Y. Qin, [Phys. Lett. B](#) **844**, 138088 (2023), [arXiv:2211.07994 \[hep-ph\]](#).
- [89] Z. Yang, Y. He, W. Chen, W.-Y. Ke, L.-G. Pang, and X.-N. Wang, [Eur. Phys. J. C](#) **83**, 652 (2023), [arXiv:2206.02393 \[nucl-th\]](#).

- [90] S. Soma, L. Wang, S. Shi, H. Stöcker, and K. Zhou, [PoS FAIRness2022](#), 055 (2023).
- [91] L. Wang, [J. Subatomic Part. Cosmol.](#) **3**, 100024 (2025), [arXiv:2501.00374 \[nucl-th\]](#).
- [92] L. Wang and J. Zhao, (2024), [arXiv:2411.16343 \[nucl-th\]](#).
- [93] L. Wang, T. Doi, T. Hatsuda, and Y. Lyu, [PoS LATTICE2024](#), 076 (2025), [arXiv:2410.03082 \[hep-lat\]](#).
- [94] T. Xu, L. Wang, L. He, K. Zhou, and Y. Jiang, [Chin. Phys. C](#) **48**, 103101 (2024), [arXiv:2405.10493 \[hep-lat\]](#).
- [95] M. Xiang, H. Yuan, L. Wang, K. Zhou, and H. G. Roskos, [IEEE Trans. Terahertz Sci. Technol.](#) **14**, 208 (2024).
- [96] K. Zhou, L. Wang, L.-G. Pang, and S. Shi, [Prog. Part. Nucl. Phys.](#) **135**, 104084 (2024), [arXiv:2303.15136 \[hep-ph\]](#).
- [97] K. Zhou, [J. Phys. Conf. Ser.](#) **2586**, 012159 (2023).
- [98] K. Hashimoto, S. Sugishita, A. Tanaka, and A. Tomiya, [Phys. Rev. D](#) **98**, 106014 (2018), [arXiv:1809.10536 \[hep-th\]](#).
- [99] K. Hashimoto, [Phys. Rev. D](#) **99**, 106017 (2019), [arXiv:1903.04951 \[hep-th\]](#).
- [100] T. Akutagawa, K. Hashimoto, and T. Sumimoto, [Phys. Rev. D](#) **102**, 026020 (2020), [arXiv:2005.02636 \[hep-th\]](#).
- [101] K. Hashimoto, K. Ohashi, and T. Sumimoto, [Phys. Rev. D](#) **105**, 106008 (2022), [arXiv:2108.08091 \[hep-th\]](#).
- [102] O.-Y. Luo, X. Chen, F.-P. Li, X.-H. Li, and K. Zhou, (2024), [arXiv:2408.03784 \[hep-ph\]](#).
- [103] M. Mansouri, K. Bitaghsir Fadafan, and X. Chen, (2024), [arXiv:2406.06285 \[hep-ph\]](#).
- [104] X. Chen and M. Huang, [Phys. Rev. D](#) **109**, L051902 (2024), [arXiv:2401.06417 \[hep-ph\]](#).
- [105] X. Chen and M. Huang, (2024), [arXiv:2405.06179 \[hep-ph\]](#).
- [106] W.-B. Chang and D.-f. Hou, [Phys. Rev. D](#) **109**, 086010 (2024), [arXiv:2403.04966 \[hep-ph\]](#).
- [107] B. Ahn, H.-S. Jeong, K.-Y. Kim, and K. Yun, [JHEP](#) **01**, 025 (2025), [arXiv:2406.07395 \[hep-th\]](#).
- [108] R.-G. Cai, S. He, L. Li, and H.-A. Zeng, (2024), [arXiv:2406.12772 \[hep-th\]](#).
- [109] K. Hashimoto, K. Matsuo, M. Murata, G. Ogiwara, and D. Takeda, [Mach. Learn. Sci. Tech.](#) **6**, 015030 (2025), [arXiv:2411.16052 \[hep-th\]](#).
- [110] B. Ahn, H.-S. Jeong, C.-W. Ji, K.-Y. Kim, and K. Yun, (2025), [arXiv:2502.10245 \[hep-th\]](#).

- [111] B. Ahn, H.-S. Jeong, K.-Y. Kim, and K. Yun, *JHEP* **03**, 141 (2024), [arXiv:2401.00939 \[hep-th\]](#).
- [112] D. Lee, H.-S. Lee, and J. Yi, (2025), [arXiv:2503.08827 \[hep-th\]](#).
- [113] O. Andreev, *Phys. Rev. D* **78**, 065007 (2008), [arXiv:0804.4756 \[hep-ph\]](#).
- [114] O. Andreev, *Phys. Rev. D* **93**, 105014 (2016), [arXiv:1511.03484 \[hep-ph\]](#).
- [115] O. Andreev, *JHEP* **05**, 173 (2021), [arXiv:2007.15466 \[hep-ph\]](#).
- [116] O. Andreev, *Phys. Rev. D* **104**, 026005 (2021), [arXiv:2101.03858 \[hep-ph\]](#).
- [117] O. Andreev, *Phys. Rev. D* **105**, 086025 (2022), [arXiv:2111.14418 \[hep-ph\]](#).
- [118] O. Andreev, *Phys. Rev. D* **107**, 026023 (2023), [arXiv:2211.12305 \[hep-ph\]](#).
- [119] O. Andreev, *Phys. Rev. D* **109**, 106001 (2024), [arXiv:2402.09026 \[hep-ph\]](#).
- [120] O. Kaczmarek and F. Zantow, *Phys. Rev. D* **71**, 114510 (2005), [arXiv:hep-lat/0503017](#).
- [121] D. Bala, O. Kaczmarek, R. Larsen, S. Mukherjee, G. Parkar, P. Petreczky, A. Rothkopf, and J. H. Weber (HotQCD), *Phys. Rev. D* **105**, 054513 (2022), [arXiv:2110.11659 \[hep-lat\]](#).
- [122] B. Chen, X. Chen, X. Li, Z.-R. Zhu, and K. Zhou, (2024), [arXiv:2404.18217 \[hep-ph\]](#).
- [123] Y. Guo, L. Qiu, R. Zhao, and M. Strickland, *Phys. Rev. D* **109**, 114025 (2024), [arXiv:2403.06739 \[hep-ph\]](#).
- [124] Q. Du, M. Du, and Y. Guo, *Phys. Rev. D* **110**, 034011 (2024), [arXiv:2402.18004 \[hep-ph\]](#).
- [125] O. Andreev, *Mod. Phys. Lett. A* **33**, 1850041 (2018), [arXiv:1707.05045 \[hep-ph\]](#).